# Analysis of the simulation

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# Thermodynamic properties

- time averages of thermodynamic quantites
  - correspond to ensemble averages (ergodic theorem)
- some quantities evaluated directly

$$U = \langle E \rangle_t$$

• fluctuations – may determine interesting properties: isochoric heat capacity:

$$C_V = \left(\frac{\partial U}{\partial T}\right)_V = \frac{\sigma_E^2}{k_B T^2} = \frac{\langle E^2 \rangle - \langle E \rangle^2}{k_B T^2}$$

- elegant way from a single simulation to heat capacity

# General note on averaging

simulated 2 MD trajectories  $\to$  two sets of 1000 values of A each perform averaging of A separately  $\to$   $\mu_1 \pm \sigma_1$  and  $\mu_2 \pm \sigma_2$ 

- how to average over the whole ensemble over 2000 values?
- $\mu = \frac{1}{2}(\mu_1 + \mu_2)$
- what about the std. deviation  $\sigma$ ?  $\sigma = \frac{1}{2}(\sigma_1 + \sigma_2)$
- hint: make use of  $\sigma^2 = \langle A^2 \rangle \langle A \rangle^2$
- solution: for each set, perform averaging of A as well as  $A^2$ , then it is safe to average the averages  $\left\langle A^2\right\rangle = \frac{1}{2}\left(\left\langle A^2\right\rangle_1 + \left\langle A^2\right\rangle_2\right)$  which leads to  $\sigma$

# Single molecule in solvent

concentrating on the dissolved molecule

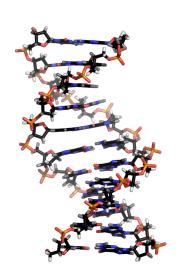
- protein, DNA,...

#### average structure

arithmetic mean of coordinates
 from snapshots along MD trajectory

$$\vec{r_i} = \frac{1}{N} \sum_{n=0}^{N} \vec{r_i}^{(n)}$$

- clear, simple, often reasonable



### Average structure

#### Possible problems:

- rotation of the entire molecule no big issue
  - RMSD fitting of every snapshot to the starting structure what is RMSD? see on the next slide...
- freely rotatable single bonds CH<sub>3</sub>
  - all 3 hydrogens collapse to a single point
  - no problem ignore hydrogens
- molecule does not oscillate around a single structure
  - several available minima of free energy
  - possibly averaging over multiple sections of trajectory

### Dynamic information

#### root mean square deviation (RMSD)

of structure in time t from a suitable reference structure  $\vec{r}^{\text{ref}}$ 

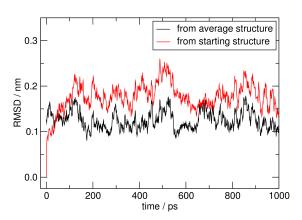
$$\mathsf{RMSD}(t) = \sqrt{rac{1}{N}\sum_{i=1}^{N}ig|ec{r_i}(t) - ec{r_i}^\mathsf{ref}ig|^2}$$

- follows the development of structure in time
- reference structure starting or average geometry
- also possible comparison with another geometry of interest DNA: A- and B-like; proteins:  $\alpha$ -helix and extended  $\beta$

RMSD fitting – finding such a translation + rotation that minimizes the RMSD from the reference structure

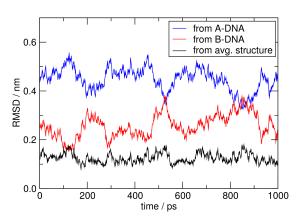
# Root mean square deviation

RMSD of non-hydrogen atoms of a DNA oligonucleotide from given geometries

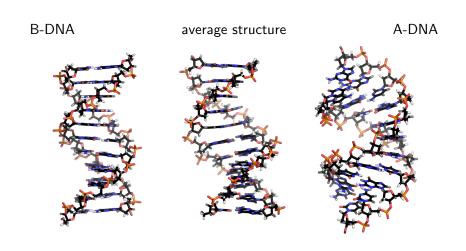


# Root mean square deviation

RMSD of non-hydrogen atoms of a DNA oligonucleotide from given geometries



# Root mean square deviation



# Magnitude of structural fluctuation

#### root mean square fluctuation (RMSF)

of position of every single atom averaged along MD trajectory (T – no. of structures in trajectory;  $\langle \vec{r_i} \rangle$  – averaged structure)

$$\mathsf{RMSF}_i = \sqrt{rac{1}{T}\sum_{t=0}^T \left| ec{r_i}(t) - \left\langle ec{r_i} 
ight
angle 
ight|^2}$$

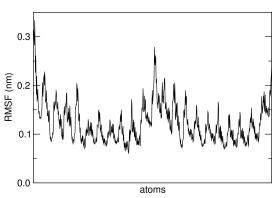
– may be converted to B-factor:

$$B_i = \frac{8}{3}\pi^2 \cdot \mathsf{RMSF}_i^2$$

- observable in diffraction experiments (X-ray...)
- contained in structure files deposited in the PDB
- comparison of simulation with X-ray may be difficult

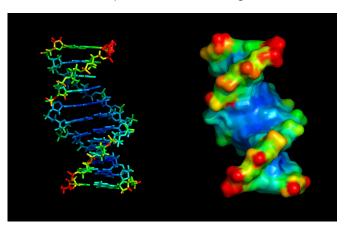
### Root mean square fluctuation

RMSF of atomic positions in DNA oligonucleotide



# Root mean square fluctuation

RMSF of atomic positions in DNA oligonucleotide



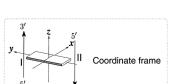
(blue < green < yellow < red)

### Structure of double-helical nucleic acids

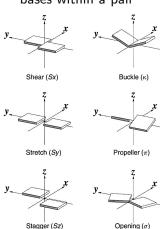


### Structure of double-helical nucleic acids

### Helical parameters

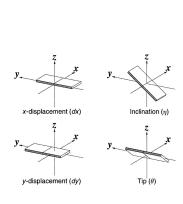


### bases within a pair

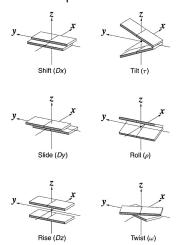


### Structure of double-helical nucleic acids

# Helical parameters pair in a helix



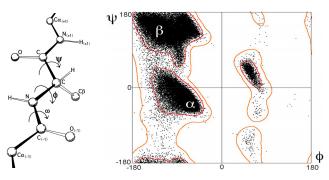
### two pairs relative



# Structure of peptides and proteins

#### Ramachandran plot

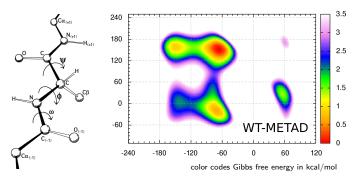
- 2D histogram of dihedrals  $\varphi$  and  $\psi$  along the backbone
- different regions correspond to various second. structures
- may be generated easily in simulation software packages



### Structure of peptides and proteins

#### Ramachandran plot

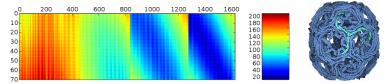
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# Structure of peptides and proteins

#### Distance matrix

- distances of amino-acid residues, represented e.g. by centers of mass or by  $C^{\alpha}$  atoms
- either time-dependent or averaged over trajectory
- bioinformatics



distance matrix between two chains (horiz. and vertical axes) shows contacts between secondary structure elements

PDB ID 1XI4, clathrin cage lattice, April 2007 Molecule of the Month

 $http://www2.warwick.ac.uk/fac/sci/moac/people/students/peter\_cock/python/protein\_contact\_map$ 

### Structure of fluids

- example pure argon or water different situation
  - many molecules, which are all equally important

#### radial distribution functions

- describe how the molecular density varies
   as a function of the distance from one particular molecule
- spherical shell of thickness  $\delta r$  at a distance r:  $\delta V \approx 4\pi r^2 \cdot \delta r$
- count the number of molecules in this shell: n
- divide by  $\delta V$  to obtain a 'local density' at distance r

### Structure of fluids

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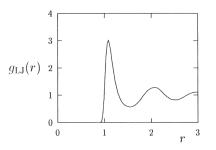
#### radial distribution functions

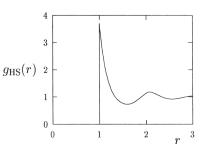
pair distribution function

$$g(r) = \frac{n/\delta V}{\rho} = \frac{n}{4\pi r^2 \cdot \delta r} \cdot \frac{1}{\rho}$$

- probability to find a molecule in distance r from ref. mol.
- division by the macroscopic density normalization

### Lennard-Jones fluid near the triple point and hard-sphere fluid





reprinted from Nezbeda, Kolafa and Kotrla 1998

- $\mathbf{g}(r)$  vanishes on short distances molecules do not intersect
- high peak van der Walls radius, closest-contact distance (even though hard spheres do not have any attraction!)
  - much more likely to find this distance in LJ or HS than in IG
- longer distances a few shallow minima and maxima, converges to unity – uniform probability as in IG

Fourier transform of g(r) – structure factor S

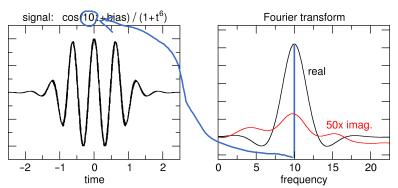
- quantifies the scattering of incoming radiation in the material
- measured in diffraction experiments (X-ray, neutron)

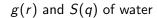
$$S(\vec{q}) = \frac{1}{N} \left\langle \sum_{i} \sum_{k} \exp\left[-i \cdot \vec{q} \cdot (\vec{r_j} - \vec{r_k})\right] \right\rangle$$

### intermission: Fourier transformation

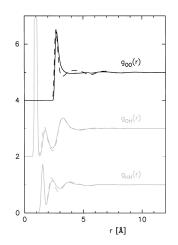
FT describes which frequencies are present in a function (of time) – decomposes f(t) into a 'sum' of periodic oscillatory functions

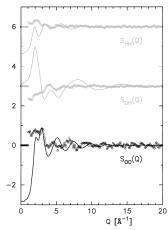
$$F(\omega) = \int_{-\infty}^{\infty} f(t) \cdot \exp\left[-\mathrm{i}\,\omega t\right] \, \mathrm{d}t \qquad \exp\left[-\mathrm{i}\,\omega t\right] = \cos\left[\omega t\right] - \mathrm{i}\sin\left[\omega t\right]$$





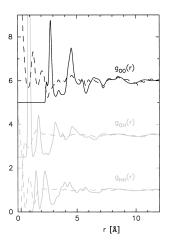
(Soper, Chemical Physics 2000)

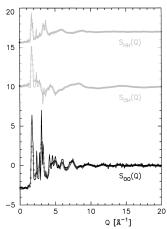




g(r) and S(q) of ice Ih at 220 K and 1 bar

(Soper, Chemical Physics 2000)





Importance – not only information about the structure calculation of thermodynamic properties possible using potential energy u(r) and force f(r) of a molecule pair corrections to the IG values of total energy and pressure (EOS!):

$$E - \frac{3}{2}Nk_{B}T = 2\pi N\rho \int_{0}^{\infty} r^{2} \cdot u(r) \cdot g(r) dr$$

$$P - \rho k_{B}T = -\frac{2\pi}{3}\rho^{2} \int_{0}^{\infty} r^{3} \cdot f(r) \cdot g(r) dr$$

(as long as pairwise additivity of forces can be assumed)

### Correlation functions

two physical quantities x(t) and y(t) may exhibit correlation

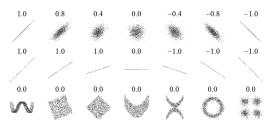
- indicates a relation of x and y, opposed to independence
- quantification several kinds of correlation functions
- Pearson correlation coefficients
  - describe linear relationship between x and y
  - quantities fluctuate around mean values  $\langle x \rangle$  and  $\langle y \rangle$
  - consider only the fluctuating part, i.e.  $x \langle x \rangle$  and  $y \langle y \rangle$
  - introduce correlation coefficient  $\rho_{xy}$

$$\rho_{xy} = \frac{\langle (x - \langle x \rangle) \cdot (y - \langle y \rangle) \rangle}{\sqrt{\langle (x - \langle x \rangle)^2 \rangle} \cdot \langle (y - \langle y \rangle)^2 \rangle} = \frac{\text{cov}(x, y)}{\sigma_x \cdot \sigma_y}$$

cov(x, y): covariance of x and y

### Correlation functions

(not necessarily linear) correlation of two quantities and the coresponding correlation coefficients



Downloaded from Wikipedia

### Correlation functions

MD – values of a quantity x as a function of time

- at some point in time, the value of x may be correlated with the value of x at an earlier time point
- described by autocorrelation function (ACF)

$$c_{x}(t) = \frac{\langle x(t) \cdot x(0) \rangle}{\langle x(0) \cdot x(0) \rangle} = \frac{\int x(t') \, x(t'+t) \, \mathrm{d}t'}{\int x^{2}(t') \, \mathrm{d}t'}$$

at two time points separated by t, normalized to takes values between -1 and 1

autocorrelation function – quantifies 'memory' of the system, or how quickly the system 'forgets' its previous state

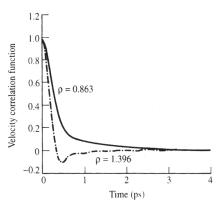
### velocity autocorrelation function

- tells how closely the velocities of atoms at time t resemble those at time 0
- usually averaged over all atoms *i* in the simulation

$$c_{\nu}(t) = rac{1}{N} \sum_{i=1}^{N} rac{\langle \vec{v_i}(t) \cdot \vec{v_i}(0) 
angle}{\langle \vec{v_i}(0) \cdot \vec{v_i}(0) 
angle}$$

• typical ACF starts at 1 in t = 0 and decreases afterwards

ACF of velocity in simulations of liquid argon (densities in g⋅cm<sup>-3</sup>)



lower  $\rho$  – gradual decay to 0

higher  $\rho$  – ACF comes faster to 0

- even becomes negative briefly
- 'cage' structure of the liquid
- one of the most interesting achievements of early simulations

Reprinted from Leach: Molecular Modelling

time needed to lose the autocorrelation whatsoever

- correlation time or relaxation time:

$$\tau_{v} = \int_{0}^{\infty} c_{v}(t) \, \mathrm{d}t$$

may help to resolve certain statistical issues: when averaging over time the properties of system, it is necessary to take uncorrelated values if the property is dynamical (related to v), we can take values of the property separated by  $\tau_{v}$ 

connection between velocity ACF and transport properties

■ Green–Kubo relation for self-diffusion coefficient *D*:

$$D = rac{1}{3} \int_0^\infty \left\langle \vec{v_i}(t) \cdot \vec{v_i}(0) 
ight
angle_i \, \mathrm{d}t$$

- interesting observable quantities
- important to be able to calculate them from MD
- there is yet another way from simulation to D
  - Einstein relation for D using mean squared displacement

$$D = \frac{1}{6} \lim_{t \to \infty} \frac{\left\langle |\vec{r_i}(t) - \vec{r_i}(0)|^2 \right\rangle_i}{t}$$

NB: Fick's laws of diffusion 
$$J=-D\frac{\partial \varphi}{\partial x}$$
,  $\frac{\partial \varphi}{\partial t}=D\frac{\partial^2 \varphi}{\partial x^2}$ 

# Autocorrelation of dipole moment

total dipole moment:

$$\vec{\mu}_{\mathsf{tot}}(t) = \sum_{i=1}^{N} \vec{\mu}_{i}(t)$$

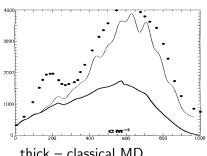
ACF of total dipole moment:

$$c_{\mu}(t) = rac{\langle ec{\mu}_{\mathsf{tot}}(t) \cdot ec{\mu}_{\mathsf{tot}}(0) 
angle}{\langle ec{\mu}_{\mathsf{tot}}(0) \cdot ec{\mu}_{\mathsf{tot}}(0) 
angle}$$

- related to the vibrational spectrum of the sample
- IR spectrum may be obtained as Fourier transform of  $c_{\mu}(t)$

# Autocorrelation of dipole moment

### IR spectra for liquid water from simulations



thick – classical MD, thin – quantum correction, black dots – experiment

B. Guillot, J. Phys. Chem. 1991

no sharp peaks at well-defined frequencies (as in gas phase)

rather – continuous bands – liquid absorbs frequencies in a broad interval

frequencies – equivalent to the rate of change of total dipole moment

covariance analysis on the atomic coordinates along MD trajectory = principal component analysis (PCA), or essential dynamics

3N-dim. covariance matrix C of atomic coordinates  $r_i \in \{x_i, y_i, z_i\}$ 

$$C_{ij} = \langle (r_i - \langle r_i \rangle) \cdot (r_j - \langle r_j \rangle) \rangle_t$$
 or  $C_{ij} = \langle \sqrt{m_i} (r_i - \langle r_i \rangle) \cdot \sqrt{m_j} (r_j - \langle r_j \rangle) \rangle_t$  (mass-weighted)

diagonalization ightarrow

eigenvalues – may be expressed as quasi-harmonic frequencies eigenvectors – principal or essential modes of motion

- analogy of normal modes of vibration
- first few largest eigenvalues, lowest frequencies
  - global, collective motions, many atoms involved

# intermission: diagonalization of a matrix

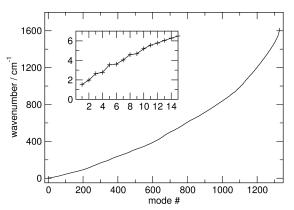
• is a process of finding eigenvectors a with corresponding eigenvalues  $\alpha$  of a square  $(n \times n)$  matrix A:

$$A \cdot a = \alpha \cdot a$$

(here: eigenvectors are column vectors)

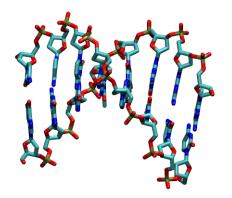
- why the name? if eigenvectors are arranged into a matrix  $\mathbf{a} = (a_1 a_2 \dots a_n)$  then  $\mathbf{a}^{-1} \cdot \mathcal{A} \cdot \mathbf{a}$  is a diagonal matrix
- lacksquare symmetric or Hermitian matrix o all eigenvalues are real
- computational cost of diagonalization:  $\mathcal{O}(n^3)$

### Double-stranded DNA oligonucleotide – lowest frequencies

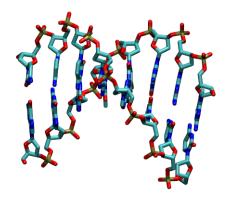


10 ns simulation of a double-helical DNA 11-nucleotide 691 atoms, of which 445 non-hydrogen ightarrow 1329 vib. modes

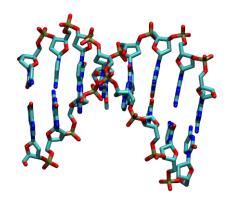
DNA octamer, eigenvector 1



DNA octamer, eigenvector 2



DNA octamer, eigenvector 3



- DNA the modes are the same as expected for a flexible rod
  - 2 bending modes around axes perpendicular to the principal axis of the DNA, and a twisting mode
- PCA gives an idea of what the modes of motion look like
  - additionally basis for thermodynamic calculations
    - vibrational frequencies may lead to configurational entropy

Binding domain of a glutamate receptor protein

