

# Molecular mechanics

classical description of molecules

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# Chemical bond

- quantum effect
- solution of Schrödinger eqn – numerically expensive  
(only small molecules can be treated)
- approximations necessary

classical treatment of nuclei

- 'field' for the moving electrons
- neglect of nuclear quantum effects (tunneling, ZPV)

Born–Oppenheimer approximation

- electrons adjust to nuclei instantaneously
- only ground state is populated

# Chemical bond

- difficult – consider delocalized electrons
- often – well localized bonding orbitals (organic molecules)
- **idea** – similar bonds have similar strength and properties  
e.g. similar C–H  $\sigma$ -orbitals  $\rightarrow$  all C–H bonds are 'similar'

How to model it?

# Chemical bond

- harmonic spring – the simplest possible function

$$E(x) = \frac{1}{2}k(x - x_0)^2$$
$$F(x) = -\frac{\partial E(x)}{\partial x} = -k(x - x_0)$$

- 2 parameters  $k$  and  $x_0$  with defined meaning
  - can be obtained from spectroscopy
- can work very well for diatomics ( $\text{H}_2$ ,  $\text{O}_2$  or  $\text{N}_2$ )

## Example – diatomic molecule

- 1D system – equiv. to 1 particle with  $m$  + spring  $k$
- Newton's second law:

$$F = ma = m \frac{d^2x}{dt^2} = -k(x - x_0)$$

- solution:

$$x(t) - x_0 = c_1 \cdot \sin \left[ \sqrt{\frac{k}{m}} t \right] + c_2 \cdot \cos \left[ \sqrt{\frac{k}{m}} t \right]$$

- eqn for harmonic motion

$$x(t) = c \cdot \sin(\omega t)$$

- relation of harmonic frequency with  $k$  and  $m$ :

$$\omega = \sqrt{\frac{k}{m}}$$

## Example – diatomic molecule

- 2nd derivative of energy equals spring constant

$$\frac{d^2 E(x)}{dx^2} = \frac{d^2 \frac{1}{2} kx^2}{dx^2} = k$$

- generally: 2nd derivative of energy determines the (harmonic) frequency of motion of atoms in molecules – **vibrations**
- **parametrization** of the simple force field is necessary
  - based on experimental data  
(equilibrium distances and vibrational frequencies)
  - or get the data from quantum chemistry

## Parametrization – general

- let us put springs in place of all covalent bonds
- do we need to parametrize  $k$  and  $x_0$  for each bond separately?  
(this would be uncool. . .)
- rather, use just several set of  $k$  and  $x_0$ ,  
one for each ‘type’ of bonding situation
- the concept of **type** is crucial

## Concept of (atom, bond... ) type

- evidence for the 'unification' of force-field parameters?

### Spectroscopy

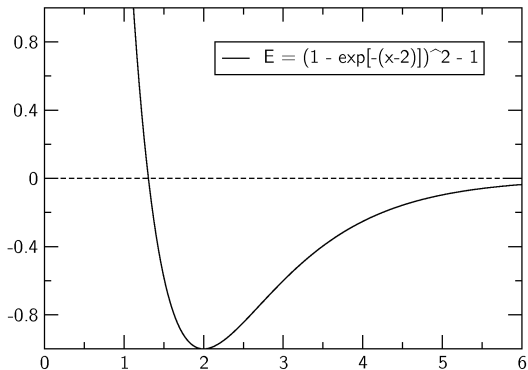
- every C–H bond: length 1.06–1.11 Å,  
frequency ca.  $3100\text{ cm}^{-1}$ , in any molecular environment

### Thermochemistry

- heat of formation – roughly additive:  
 $\text{CH}_4 \cong 4\text{ C-H}$   
 $\text{C}_2\text{H}_6 \cong 6\text{ C-H} + \text{C-C}$
- connection to well depth in Morse's potential

# Concept of (atom, bond...) type

## Morse's potential



# Concept of (atom, bond... ) type

Bonds between atoms to be modeled by universal functions  
How to identify the **atom types**? – chemical ideas

i) hybridization

- difference between  $sp^3$  carbon (4 bonds) and  $sp^2$  C (3 bonds)  
– these will be the different atom types
- different functions for bonds of types C–C, C=C and C≡C
- determine the parameters ( $k$ ,  $x_0$ ) with some selected molecules, typical for the binding situation
- example: use  $C_2H_6$ ,  $C_2H_4$ ,  $C_2H_2$  and benzene for  $k$ ,  $x_0$

# Concept of (atom, bond...) type

Bonds between atoms to be modeled by universal functions

How to identify the **atom types**? – chemical ideas

ii) polarity

- an atom bonded to electronegative atom – electron deficient  
→ affects its bonding to other atoms
- example:  $\text{O}=\text{CH}-\text{C}\dots$   
the C–C bond is affected and needs to be parametrized  
differently from apolar C–C  
→ an atom type for carbonyl C introduced

Biomolecular force fields

– usually 20 types for C, 10 for N and 5 for O and H

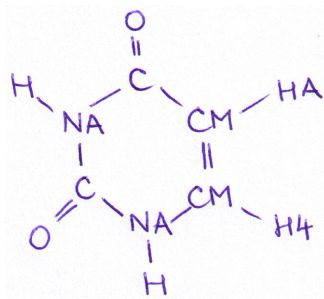
## Concept of (atom, bond...) type

AMBER force field for biomolecules:

- C - sp<sup>2</sup> C carbonyl group
- CA - sp<sup>2</sup> C pure aromatic (benzene)
- CB - sp<sup>2</sup> aromatic C, 5&6 membered ring junction
- CC - sp<sup>2</sup> aromatic C, 5 memb. ring HIS
- CK - sp<sup>2</sup> C 5 memb.ring in purines
- CM - sp<sup>2</sup> C pyrimidines in pos. 5 & 6
- CN - sp<sup>2</sup> C aromatic 5&6 memb.ring junct.(TRP)
- CQ - sp<sup>2</sup> C in 5 mem.ring of purines between 2 N
- CR - sp<sup>2</sup> arom as CQ but in HIS
- CT - **sp<sup>3</sup>** aliphatic C
- CV - sp<sup>2</sup> arom. 5 memb.ring w/1 N and 1 H (HIS)
- CW - sp<sup>2</sup> arom. 5 memb.ring w/1 N-H and 1 H (HIS)
- C\* - sp<sup>2</sup> arom. 5 memb.ring w/1 subst. (TRP)

## Concept of (atom, bond...) type

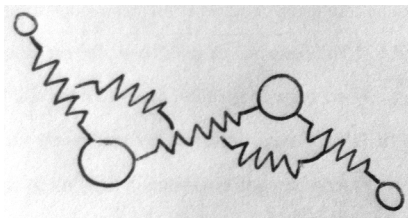
AMBER atom types in a molecule of uracil



# Interactions between atoms

## Bonded

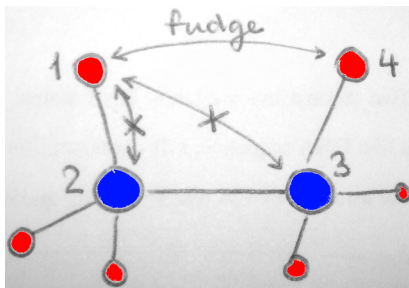
- mediated by, and resulting directly from covalent bonds
- usually described by harmonic springs between atoms, and bond and dihedral angles in a similar way
- cover all of the quantum-mechanical phenomena like exchange and correlation between two atoms using an effective potential function



# Interactions between atoms

## Non-bonded

- longer-range interactions
  - charge–charge Coulomb and van der Waals (vdW)
- between molecules and distant parts of one molecule
- **excluded** between bonded atoms, up to 4th neighbor
- 1–4 interactions – scaled down ('fudge' factor)



# Coulomb parameters

3 contributions (quantum-chemical picture)

- nucleus–nucleus

$$\frac{1}{2} \sum_{ij} \frac{Z_i \cdot Z_j}{R_{ij}}$$

- nucleus–electron with approx. point charges  $\rho(r) \rightarrow \sum_j Q_j$

$$-\sum_i \int \frac{Z_i \cdot \rho(r)}{|R_i - r|} dr \rightarrow -\sum_{ij} \frac{Z_i \cdot Q_j}{R_{ij}}$$

- (Hartree) electron–electron with point charges

$$\frac{1}{2} \iint \frac{\rho(r) \cdot \rho(r')}{|r - r'|} \rightarrow \frac{1}{2} \sum_{ij} \frac{Q_i \cdot Q_j}{R_{ij}}$$

# Coulomb parameters

3 contributions (quantum-chemical picture)

- nucleus–nucleus
- nucleus–electron with approx. point charges  $\rho(r) \rightarrow \sum_j Q_j$
- (Hartree) electron–electron with point charges

define an **effective atomic charge** as  $q_i = -Q_i + Z_i$   
 then

$$E_{QQ} = \frac{1}{4\pi\epsilon_0} \cdot \sum_i \sum_{j>i} \frac{q_i \cdot q_j}{R_{ij}}$$

## Coulomb parameters

effective atomic charge needs to be defined for every atom  
(rather than atom type – this would be too crude)

- from quantum-chemical calculations – straightforward
- do calculations of typical (bio)molecular fragments and get charges from there
- fragments
  - amino acid residues and peptide bonds for proteins
  - nucleobases, sugars and phosphate groups for DNA/RNA
- there are 2 issues...

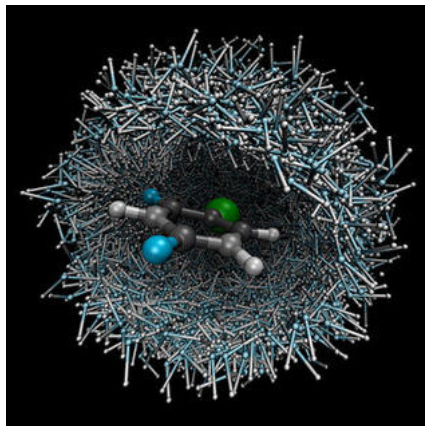
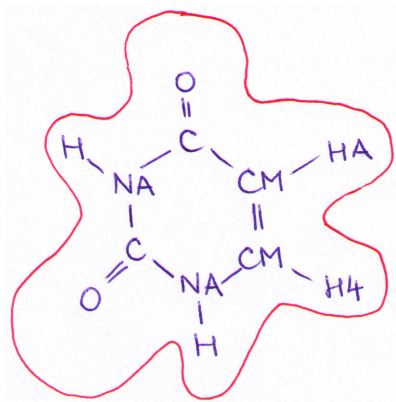
# Coulomb parameters – issue 1

atomic charges – no unique definition in quantum chemistry  
(Mulliken, or may it be NBO or AIM today?)

- several schemes proposed – hard to say, which is the right
- the simplest (Mulliken charges) do not work
- popular – **potential-derived** charges:
  - 1 calculate the electron density in the molecule
  - 2 get **electrostatic potential** at surface of the molecule
  - 3 **fit** point electric charges on atoms to reproduce the ESP

# Coulomb parameters – issue 1

surface of the uracil molecule



## Coulomb parameters – issue 2

atom charges are calculated for an isolated molecule (gas phase)

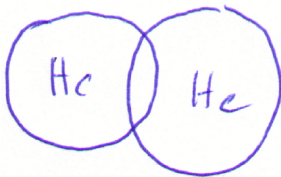
- electrostatic situation in aqueous/polar solution differs
  - molecules are more polar
- example: H<sub>2</sub>O molecule in the gas phase: 1.8 D  
in water: 2.4 D
- a good idea – to set the charges to larger values
- popular trick – use a small basis sets in QCh calculations
  - then, charges and dipole moments are overestimated
  - good for force fields w/o polarizability and charge transfer
- another possibility – **polarizable force field**

atomic polarizability  $\alpha_i$  is assigned to every atom  $i$   
external field induces atomic dipole  $\vec{\mu}_i = \vec{\alpha}_i \cdot \vec{E}$

# van der Waals interaction

## Pauli repulsion

- electrons with the same spin avoid spatial overlap
- purely quantum-mechanical phenomenon (exchange effect)
- example: interaction of neutral closed-shell systems – He<sub>2</sub>



- electron densities start to overlap → energy grows steeply
- modeling:

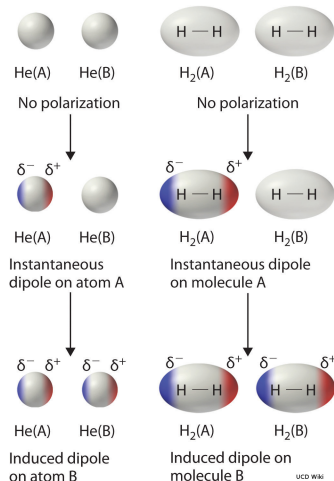
$$E_{\text{ex}} = \exp[a - b \cdot R_{ij}]$$

$$E_{\text{ex}} = \left( \frac{\sigma}{R_{ij}} \right)^{12}$$

# van der Waals interaction

## dispersion due to correlation

- correlation – between electrons irrespective of spin, retained on longer distances
- fluctuation of density
  - instantaneous dipole
  - induced dipole on another atom
  - interaction
- orientation of dipoles is correlated
  - **attractive** interaction
- $R^{-6}$ -dependence, proportional to polarizabilities



note: dispersion is difficult to capture with quantum chemistry

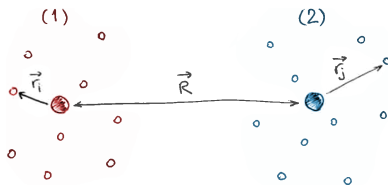
# van der Waals interaction

dispersion due to correlation



# van der Waals interaction

dispersion – quantification (1)



$i, j$  are electrons with charges  $e$ ;  $\vec{R} \parallel \vec{z}$ ;  $\vec{r}_i = (x_i, y_i, z_i)$ ;  $X_1 = \sum_i x_i$   
interaction of charge distributions (leading Taylor term):

$$V = \frac{1}{R^3} \sum_{i,j} e_i e_j (x_i x_j + y_i y_j - 2z_i z_j) = \frac{e^2}{R^3} (X_1 X_2 + Y_1 Y_2 - 2Z_1 Z_2)$$

**perturbation theory** – interaction of dipoles is the perturbation

# van der Waals interaction

dispersion – quantification (2)

no permanent dipole  $\rightarrow$  first-order energy vanishes:

$$E^{(1)} = \langle \Psi^{(0)} | V | \Psi^{(0)} \rangle = 0$$

second-order energy w/ unperturbed states  $|\Psi_m^{(0)}\rangle$  & energies  $E_m^{(0)}$ :

$$E^{(2)} = \sum_{m \neq n} \frac{|\langle \Psi_m^{(0)} | V | \Psi_n^{(0)} \rangle|^2}{E_n^{(0)} - E_m^{(0)}}$$

recall  $V = \frac{e^2}{R^3}(X_1 X_2 + Y_1 Y_2 - 2Z_1 Z_2)$

and write  $X_{m0} = \langle \Psi_m^{(0)} | \frac{eX}{R^3} | \Psi_0^{(0)} \rangle$

# van der Waals interaction

dispersion – quantification (3)

$$E^{(2)} = \frac{e^4}{R^6} \sum_{m \neq n} \frac{|X_{m0}^1 X_{n0}^2 + Y_{m0}^1 Y_{n0}^2 - 2Z_{m0}^1 Z_{n0}^2|^2}{E_{10}^{(0)} - E_{1m}^{(0)} + E_{20}^{(0)} - E_{2n}^{(0)}}$$

$R^{-6}$  arises from the matrix element being squared  
 simplification – averaging over all orientations of clusters 1 and 2,  
 and also considering spherically symmetric charge distributions:

$$E^{(2)} = \frac{2}{3} \frac{e^4}{R^6} \sum_{m \neq n} \frac{|R_{m0}^1|^2 |R_{n0}^2|^2}{E_{10}^{(0)} - E_{1m}^{(0)} + E_{20}^{(0)} - E_{2n}^{(0)}}$$

# van der Waals interaction

dispersion – quantification (4)

consider polarizability  $\alpha(0)$  in a static field:

$$\alpha(0) = \frac{2}{3} e^2 \sum_j \frac{|R_{j0}|^2}{E_j - E_0} \approx \frac{2}{3} \frac{e^2}{I} \sum_j |R_{j0}|^2$$

with ionization energy  $I$  approximating excitation energy  $E_j - E_0$

$$\begin{aligned} E^{(2)} &\approx -\frac{2}{3} \frac{e^4}{R^6} \frac{1}{I_1 + I_2} \sum_{m \neq n} |R_{m0}^1|^2 |R_{n0}^2|^2 \\ &\approx -\frac{2}{3} \frac{I_1 I_2}{I_1 + I_2} \frac{\alpha_1(0) \alpha_2(0)}{R^6} = -\frac{C_6}{R^6} \end{aligned}$$

Note – higher-order terms have been neglected

# van der Waals interaction

most common function: **Lennard-Jones 12-6** potential

$$V(r) = 4\varepsilon \left( \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right)$$

- 2 parameters –  $\sigma$  and  $\varepsilon$
- minimum at  $r_m = 2^{1/6}\sigma$  and well depth  $\varepsilon$

repulsive:  $\exp[-r]$  sometimes better than  $r^{-12}$

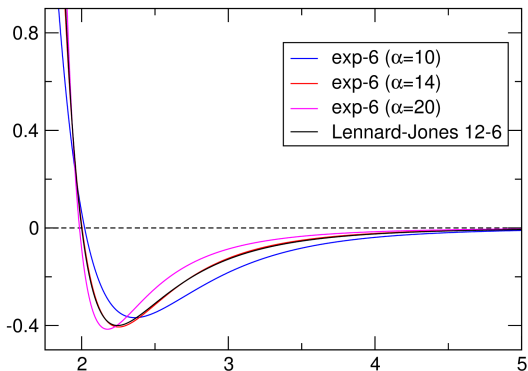
hskip 5 mm → **exp-6** potential a.k.a. Buckingham potential

$$V(r) = \frac{\varepsilon}{1 - 6/\alpha} \cdot \left( \frac{6}{\alpha} \cdot \exp \left[ \alpha \left( 1 - \frac{r}{\sigma} \right) \right] - \left( \frac{\sigma}{r} \right)^6 \right)$$

- $\alpha$  controls the steepness of the repulsive branch
- may be a better choice for phase transitions
- e.g. MM water would not freeze below 0 °C with LJ 12-6

# van der Waals interaction

Lennard-Jones 12-6 potential with  $\varepsilon = 0.4$  and  $\sigma = 2$ , and several similar exp-6 potentials with various values of  $\alpha$



# van der Waals interaction

heteronuclear interactions A–B

– mixing rules for the parameters

$$\sigma_{AB} = \frac{\sigma_{AA} + \sigma_{BB}}{2}$$
$$\varepsilon_{AB} = \sqrt{\varepsilon_{AA} \cdot \varepsilon_{BB}}$$

parametrization

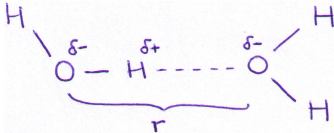
- challenging task in general
- fitting of params to experimental / quantum-chem. data
- e.g. relation to density and heat of vaporization
  - obvious in organic liquids – major interaction

# Hydrogen bonding

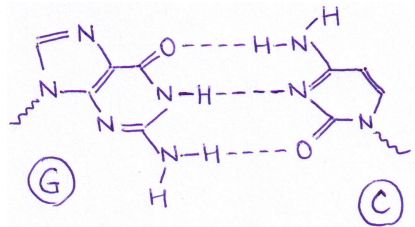
- attractive interaction between an H bonded to an electronegative atom (O, N) and another O or N
- examples: H<sub>2</sub>O dimer, nucleobase pair
- crucial interaction in biomolecules
- interplay of several kinds of interactions
- early force fields – special potential functions for H-bonding
- typical binding energies: 20 kJ/mol
  - higher for strongly polarized or even charged molecules
  - or if there are several H-bonds (nucleobase pairs)

# Hydrogen bonding

H<sub>2</sub>O dimer



guanine:cytosine base pair

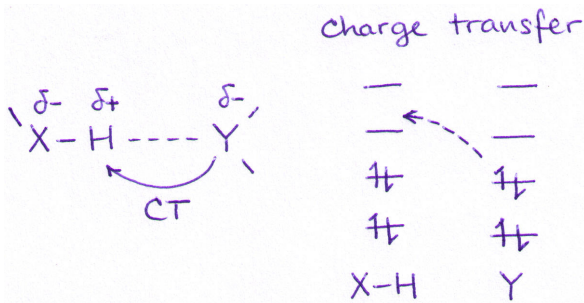


# Hydrogen bonding

- Coulomb interaction is dominant
- vdW interaction
  - may become important, especially in weakly bound systems
  - crucial e.g. for angular dependence in  $\text{H}_2\text{CO} \dots \text{H}_2\text{O}$  etc.
- modern force fields – no special treatment of H-bonds
- **charge transfer** contribution
  - cannot be covered by force fields due to constant charges
  - may be included in other terms effectively

# Hydrogen bonding

Charge transfer between the charge donor (Y) and acceptor (X-H)  
 (or proton donor X-H and proton acceptor Y)



# Parametrization of bonded interactions

## Bonds

- usually represented by Morse's potential well

$$E(r) = D (1 - \exp[-\alpha(r - r_0)])^2$$

– not necessary, possibly inefficient, hardly used

- way to an approximated function – Taylor expansion

$$E(r) = E(r_0) + \frac{dE}{dr}(r_0) \cdot (r - r_0) + \frac{1}{2} \frac{d^2E}{dr^2}(r_0) \cdot (r - r_0)^2 + \dots$$

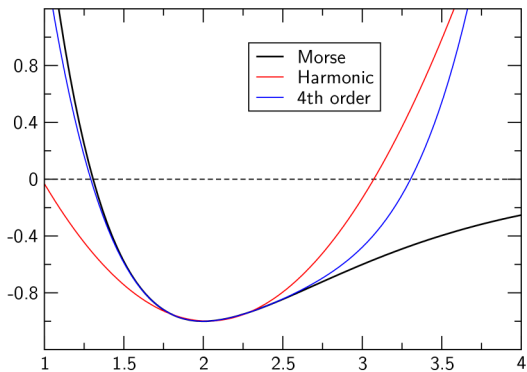
- most force fields – harmonic approximation  
 (Taylor expansion cut after 2nd-order term)

$$E(r) = \frac{1}{2} k (r - r_0)^2$$

- approximation of the true potential with a quadratic function  
 – works in a narrow range of distances

# Parametrization of bonded interactions

Comparison of the harmonic and quartic functions with Morse's



- for accurate vibration frequencies – 4th-order terms can be important to describe the curvature

# Parametrization of bonded interactions

## Angles

- harmonic approximation for the angle deformation

$$E_{\text{bend}}(\vartheta) = \frac{1}{2}k_{\vartheta}(\vartheta - \vartheta_0)^2$$

- two parameters needed – equilibrium angle and force constant
- from experiment (vib-rot spectra) or quantum chemistry

# Parametrization of bonded interactions

## Dihedral angles

- describe the rotation around covalent bonds
- defined by 4 atoms
- potential energy – periodic function of the dihedral angle:

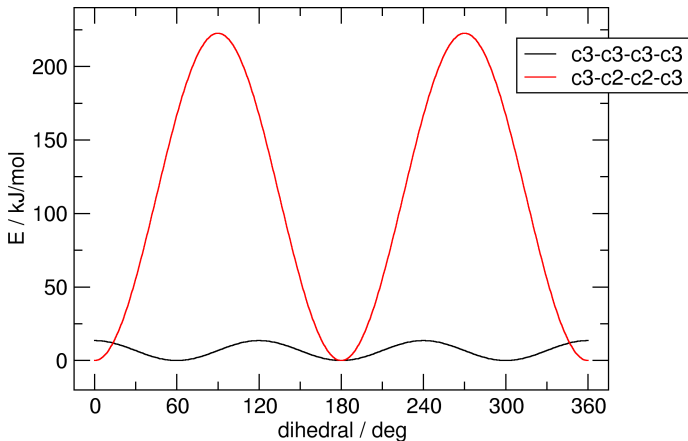
$$E(\omega) = \sum_{n=1,2,3,4,6} V_n \cos [n \cdot \omega - \gamma_n]$$

- $V_n$  – amplitude (barrier),  $n$  – periodicity,  $\gamma$  – phase shift
- Ryckaert-Bellemans function – often used for alkanes:

$$E(\Omega) = \sum_{n=0}^5 C_n \cos^n \Omega \quad (\Omega = \omega - 180^\circ)$$

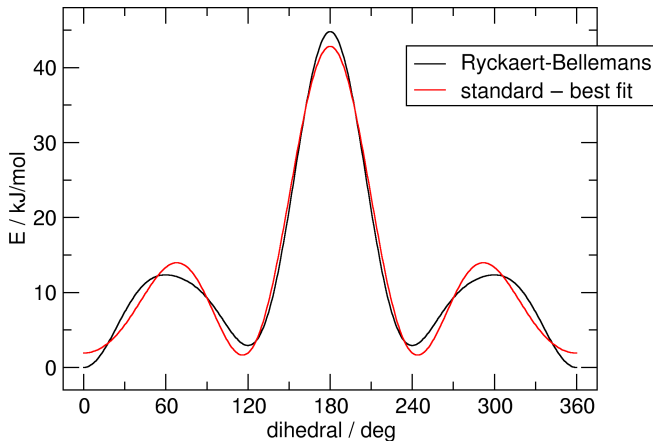
# Parametrization of bonded interactions

Dihedral angles – example: C–C single and C=C double bonds



# Parametrization of bonded interactions

Dihedral angles – comparison of RB dihedral with standard form



# The complete equation

Adding up all contributions, the total energy of a typical biomolecular force field reads

$$\begin{aligned}
 E(R^N) = & \\
 = & \frac{1}{2} \sum_i k_i (r_i - r_i^0)^2 + \frac{1}{2} \sum_j k_j^\vartheta (\vartheta_j - \vartheta_j^0)^2 + \frac{1}{2} \sum_n V_n \cdot \cos[n\omega - \gamma_n] \\
 + & \sum_i^N \sum_{j=i+1}^N \left\{ 4\epsilon_{ij} \left( \left( \frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left( \frac{\sigma_{ij}}{r_{ij}} \right)^6 \right) + \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{r_{ij}} \right\}
 \end{aligned}$$

# Forces

force acting on atom  $i$

– derivative of energy with respect to the coordinates of atom  $i$

$$\begin{aligned}\vec{F}_i &= -\nabla_i V \\ F_i^x &= -\frac{\partial V}{\partial x_i}\end{aligned}$$

derivatives of all of the terms in the force field  
 are obtained in an analytical form easily

# Forces

most terms depend on distances of atoms

$$V = V(r_{ij}) + V(r_{ik}) + V(r_{il}) + \dots$$

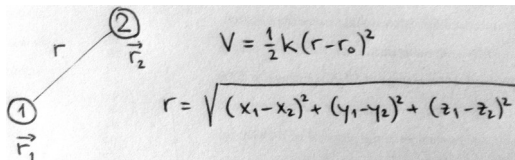
and thus the force

$$\begin{aligned} F_i^x &= -\frac{\partial V}{\partial x_i} = \\ &= -\frac{\partial V(r_{ij})}{\partial x_i} - \frac{\partial V(r_{ik})}{\partial x_i} - \frac{\partial V(r_{il})}{\partial x_i} - \dots = \\ &= -\frac{\partial V(r_{ij})}{\partial r_{ij}} \frac{\partial r_{ij}}{\partial x_i} - \frac{\partial V(r_{ik})}{\partial r_{ik}} \frac{\partial r_{ik}}{\partial x_i} - \frac{\partial V(r_{il})}{\partial r_{il}} \frac{\partial r_{il}}{\partial x_i} - \dots \end{aligned}$$

(more complex derivatives for angle/dihedral contributions)

$$\frac{\partial V(\angle ijk)}{\partial \angle ijk} \frac{\partial \angle ijk}{\partial x_i}, \quad \frac{\partial V(\angle ijkl)}{\partial \angle ijkl} \frac{\partial \angle ijkl}{\partial x_i}$$

## Forces – example



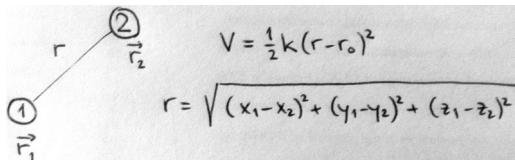
$$V(r_{12}) = \frac{1}{2}k(r_{12} - r_0)^2$$

$$\frac{\partial V(r_{12})}{\partial r_{12}} = k(r_{12} - r_0)$$

$$r_{12} = \sqrt{(x_1 - x_2)^2 + (y_1 - y_2)^2 + (z_1 - z_2)^2}$$

$$\frac{\partial r_{12}}{\partial x_1} = \frac{1}{2r_{12}} \cdot \frac{\partial (x_1 - x_2)^2}{\partial x_1} = \frac{1}{2r_{12}} \cdot 2(x_1 - x_2) = \frac{x_1 - x_2}{r_{12}}$$

## Forces – example



$$\frac{\partial V}{\partial x_1} = \frac{\partial V(r_{12})}{\partial r_{12}} \frac{\partial r_{12}}{\partial x_1} = k(r_{12} - r_0) \cdot \frac{x_1 - x_2}{r_{12}}$$

$$\vec{F}_1 = -\nabla_1 V = -k(r_{12} - r_0) \cdot \frac{\vec{r}_{12}}{r_{12}} \quad \text{with } \vec{r}_{12} = \vec{r}_1 - \vec{r}_2$$